



# AS3020: Aerospace Structures

## Module 4: Bending of Beam-Like Structures

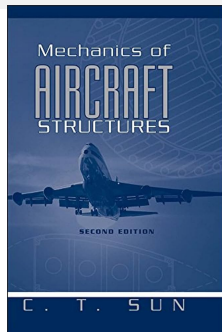
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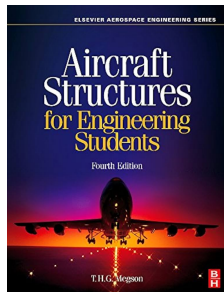
September 13, 2024

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*Chapters 4-5 in Sun [1]*

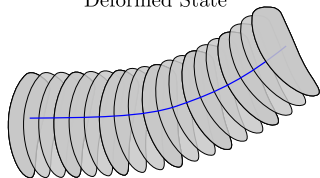


*Chapters 16-20 in Megson [2]*

# 1. Unsymmetrical Bending



Deformed State



- Displacement Field

$$u_1 = -X_2\theta_3 + X_3\theta_2, \quad u_2 = v, \quad u_3 = w.$$

- Zero shear  $\implies \theta_3 = v'$ ,  $\theta_2 = -w'$
- Direct stress

$$\begin{aligned} \sigma_{11} &= E_y [X_3 \quad -X_2] \begin{bmatrix} \theta_2' \\ \theta_3' \end{bmatrix} \\ &= \frac{[X_3 \quad -X_2]}{I_{22}I_{33} - I_{23}^2} \begin{bmatrix} I_{33}M_2 + I_{23}M_3 \\ I_{23}M_2 + I_{22}M_3 \end{bmatrix} \end{aligned}$$

- Equilibrium Considerations:

$$M_{2,1} = V_3, \quad V_{2,1} + F_2 = 0$$

$$M_{3,1} = -V_2, \quad V_{3,1} + F_3 = 0.$$

## 2. Shear Stress and Flow in Sections

- If shear strain is assumed zero, can we still have shear stress?
- We posit:  $\gamma_{12} = 0$ ,  $\gamma_{13} = 0$ ,  $\gamma_{23} = 0$ . As point quantities, the shear stresses may still be small ( $\tau_{12} = G\sigma_{12}$ ).
- But the **integral quantities** are taken to be finite:

$$\int \sigma_{12} dA = V_2, \quad \int \sigma_{13} dA = V_3,$$

$$\int \sigma_{12} dX_3 = q_2, \quad \int \sigma_{13} dX_2 = q_3.$$

- Invoking plane stress assumption at the section, the governing equation is,

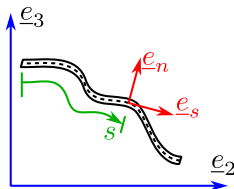
$$\sigma_{11,1} + \sigma_{1s,s} = 0.$$

- Integrating the above from  $s = 0$  to  $s$ , we get the **Shear flow formula**:

$$q(s) - q_0 = - \frac{\left[ \int_0^s tX_3 ds \quad - \int_0^s tX_2 ds \right]}{I_{22}I_{33} - I_{23}^2} \begin{bmatrix} I_{33}V_3 - I_{23}V_2 \\ I_{23}V_3 - I_{22}V_2 \end{bmatrix}$$

## 2. Shear Stress and Flow in Sections

Thin Section: Plane Stress Assumption



- We define the above section-local coordinate system and transform the elasticity equations to  $\sigma_{11,1} + \sigma_{1n,n} + \sigma_{1s,s} = 0$ . Applying **plane stress** assumption (for thin sections) drops the  $\sigma_{1n}$  term, leading to:

$$\sigma_{11,1} + \sigma_{1s,s} = 0 \implies t\sigma_{11,1} + q_{,s} = 0,$$

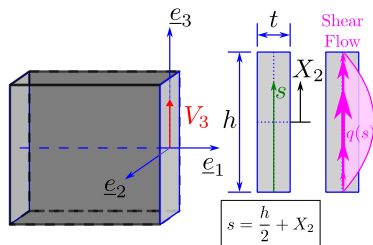
where we have integrated along the  $\underline{e}_n$  direction once.

- Following through with the integral along  $\underline{e}_s$ , this leads to the **shear flow formula**

$$q(s) - q_0 = - \frac{\left[ \int_0^s tX_3 ds \quad - \int_0^s tX_2 ds \right]}{I_{22}I_{33} - I_{23}^2} \begin{bmatrix} I_{33}V_3 - I_{23}V_2 \\ I_{23}V_3 - I_{22}V_2 \end{bmatrix}$$

## 2. Shear Stress and Flow in Sections

- Consider the rectangular section with height  $h$  and thickness  $t$ :



$$\begin{aligned}
 q(s) &= -\frac{V_3}{I_{22}} \int_0^s t X_3 ds = -\frac{tV_3}{I_{22}} \int_{-\frac{h}{2}}^{X_3} X_3 dX_2 \\
 &= -\frac{tV_3}{2I_{22}} \left( X_2^2 - \frac{h^2}{4} \right)
 \end{aligned}$$

- Remember that  $V_3$  is NOT any externally **applied force**. It is merely **the resultant of all the shear stresses in the section**.
- So  $V_3$  and  $q(s)$  point in the same direction in this example. It is incorrect to think that  $q(s)$  is balancing out  $V_3$ .

## 2. Shear Stress and Flow in Sections

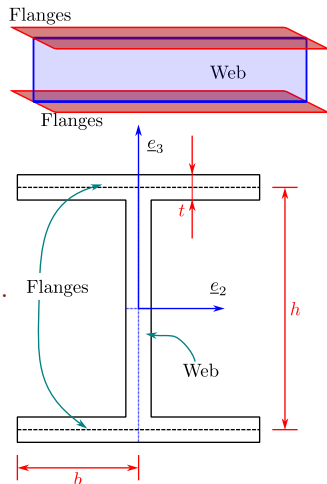
### The "I" section

- Consider the shear distribution through an I-section as shown here
- The shear distribution looks like it is “flowing”, with more “flow” occurring in the thin vertical web and less in the flanges.
- The second moment of area  $I_{22}$  sums up as,

$$I_{22} = \underbrace{\frac{h^3 t}{12}}_{\text{web}} + 2 \times \underbrace{\left( \frac{2bt^3}{12} + 2bt \times \frac{h^2}{4} \right)}_{\text{flange}} \approx \left( \frac{h^3}{12} + bh^2 \right) t.$$

- $I_{33}$  sums up as,

$$I_{33} = \underbrace{\frac{ht^3}{12}}_{\text{web}} + 2 \times \underbrace{\left( \frac{2b^3 t}{3} \right)}_{\text{flange}} \approx \underbrace{\frac{4b^3 t}{3}}_{\approx 0 \text{ for small } b}.$$

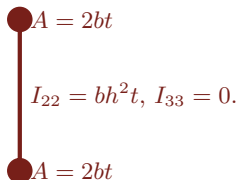


## 2. Shear Stress and Flow in Sections

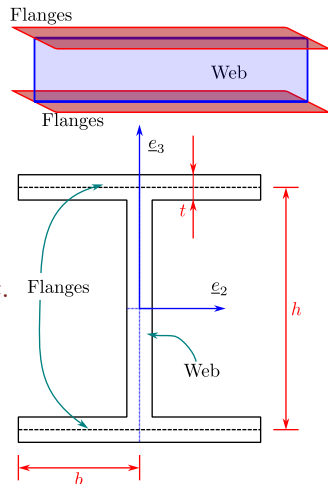
The “I” section

### Idealization

- Both  $I_{22}$  and  $I_{33}$  are dominated by flange contributions, implying that bending is supported primarily by the flanges.
- This motivates the following idealization for the I-section:



- The lumped area elements denoted ● are sometimes referred to as “**Booms**” in the section.
- Thickness in the web (denoted  $t$ ) is taken to be zero for bending-stress calculations.

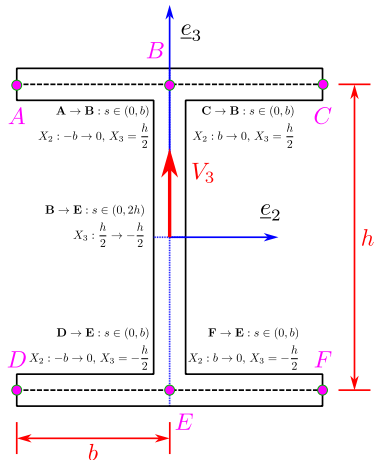




## 2. Shear Stress and Flow in Sections

### The "I" section

- Let us consider the case with  $V_2 = 0$ ,  $V_3 \neq 0$ .
- The segments  $A \rightarrow B$ ,  $C \rightarrow B$ ,  $D \rightarrow E$ , and  $F \rightarrow E$  are exposed in their free ends, simplifying the shear flow integral ( $q = 0$  at free ends).

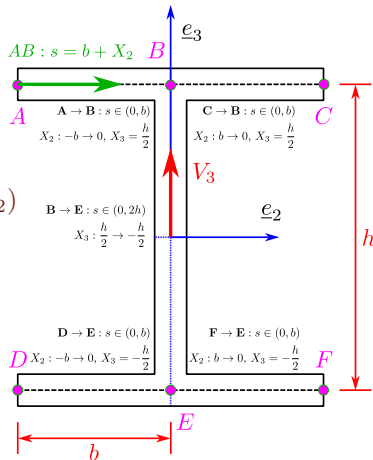


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$$\mathbf{A} \rightarrow \mathbf{B} : q_{AB}(s) \equiv q_{AB}(X_2) = -\frac{tV_3}{I_{22}} \int_0^s X_3 ds = -\frac{htV_3}{2I_{22}}(b + X_2)$$



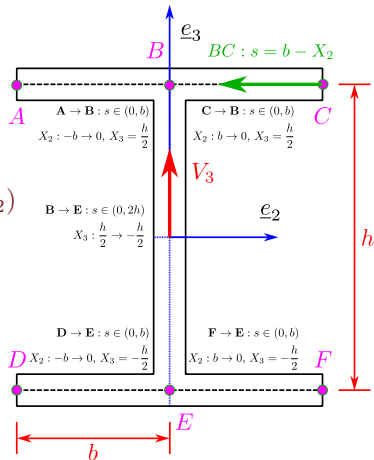
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$$\mathbf{C} \rightarrow \mathbf{B} : q_{CB}(X_2) = -\frac{htV_3}{2I_{22}}(b - X_2)$$



## 2. Shear Stress and Flow in Sections

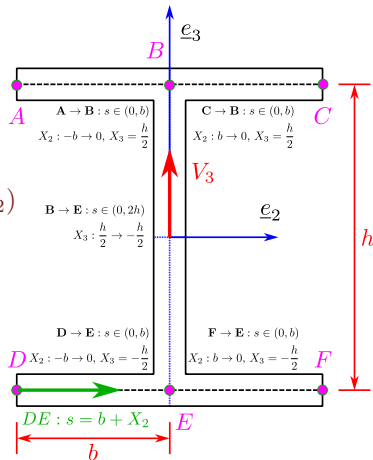
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$$\mathbf{D} \rightarrow \mathbf{E} : q_{DE}(X_2) = \frac{htV_3}{2I_{22}}(b + X_2)$$



## 2. Shear Stress and Flow in Sections

### The "I" section

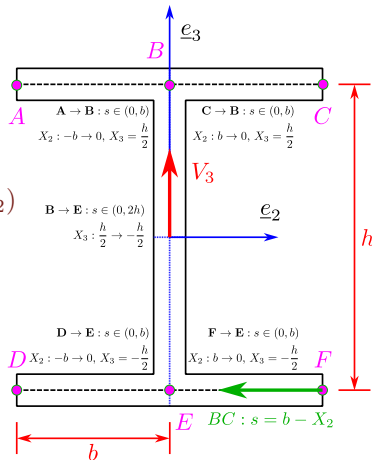
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$$\mathbf{D} \rightarrow \mathbf{E} : q_{DE}(X_2) = \frac{htV_3}{2I_{22}}(b + X_2)$$

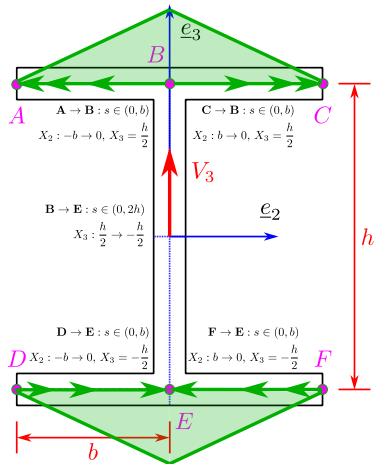
$$\mathbf{F} \rightarrow \mathbf{E} : q_{FE}(X_2) = \frac{htV_3}{2I_{22}}(b - X_2)$$



## 2. Shear Stress and Flow in Sections

### The "I" section

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- In summary we have linear relationships at the flanges.

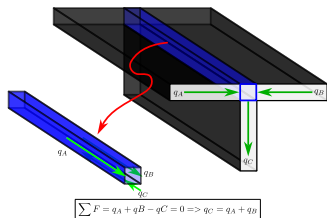


## 2. Shear Stress and Flow in Sections

### The "I" section

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- In summary we have linear relationships at the flanges.
- Before looking at the web ( $B \rightarrow E$ ), we have to observe the balance at the "T" junction.

### Balance at the T-junction



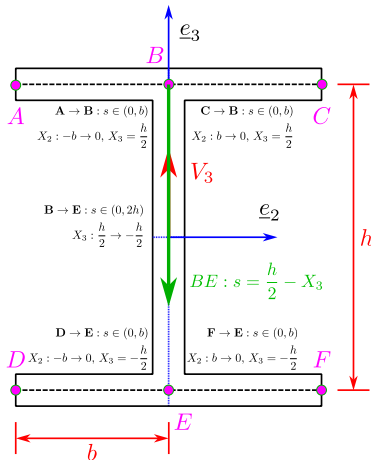
## 2. Shear Stress and Flow in Sections

### The "I" section

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- The segments  $A \rightarrow B$ ,  $C \rightarrow B$ ,  $D \rightarrow E$ , and  $F \rightarrow E$  are exposed in their free ends, simplifying the shear flow integral ( $q = 0$  at free ends).
- On  $B \rightarrow E$ , we have  

$$q_{BE}(0) = q_{AB}(b) + q_{CB}(b) = -\frac{bhtV_3}{I_{22}}.$$
- The integration evaluates as,

$$\begin{aligned} q_{BE}(X_3) &= -\frac{bhtV_3}{I_{22}} + \frac{tV_3}{2I_{22}}(X_3^2 - \frac{h^2}{4}) \\ &= -\frac{htV_3}{I_{22}}(b + \frac{h}{8}) + \frac{tV_3}{2I_{22}}X_3^2. \end{aligned}$$





## 2. Shear Stress and Flow in Sections

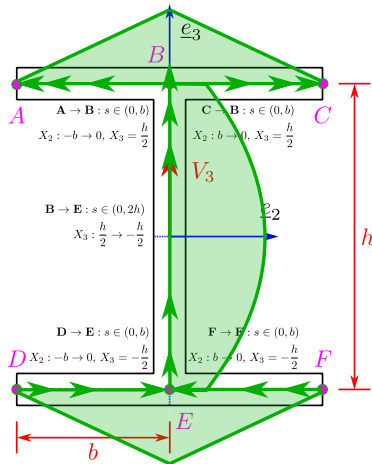
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- We now have the complete shear flow in the section.



## 2. Shear Stress and Flow in Sections

The “I” section: Order of Magnitude Analysis

- Let us consider the “total” shear forces experienced by each member.
- Flange AB**

$$V_{AB} = -\frac{htV_3}{2I_{22}} \int_{-b}^0 (b + X_2) dX_2 = -\frac{b^2htV_3}{4I_{22}}$$

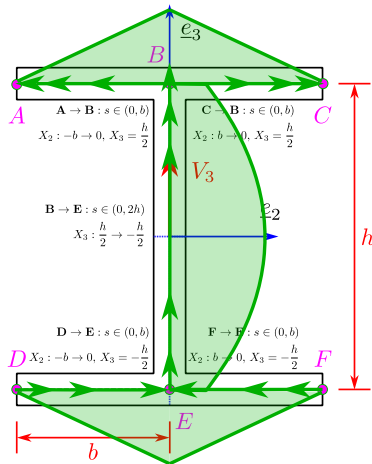
- Web BE**

$$V_{BE} = \frac{h^2(h + 12b)tV_3}{12I_{22}} = V_3$$

- For  $b = \frac{h}{2}$ , we have,

$$V_{AB} = -\frac{h^3tV_3}{16I_{22}} \approx -\frac{V_3}{8}$$

$$V_{BE} = V_3$$

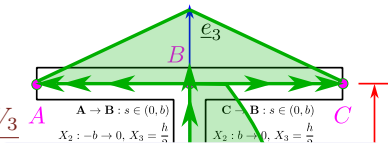


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The “I” section: Order of Magnitude Analysis

- Let us consider the “total” shear forces experienced by each member.
- Flange AB**

$$V_{AB} = -\frac{htV_3}{2I_{22}} \int_{-b}^0 (b + X_2) dX_2 = -\frac{b^2htV_3}{4I_2}$$



**Idealization**

Since  $V_{AB} \ll V_{BE}$ , we understand that the **web is primarily responsible for restoring shear loads**, with negligible contributions from the flanges.

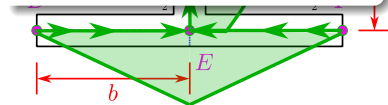
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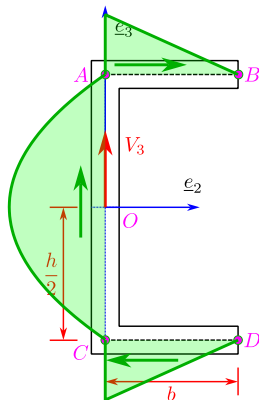
$$V_{AB} = -\frac{h^3tV_3}{16I_{22}} \approx -\frac{V_3}{8}$$

$$V_{BE} = V_3$$



## 2. Shear Stress and Flow in Sections

- Shear Center is the point of shear load such that  $\sum m_1 = 0$ . Although in a lot of symmetric sections this is coincident with the centroid, **this is NOT always the case.**
- Consider the “C” section beam:



$$I_{22} \approx \frac{(h^3 + 6bh^2)t}{12} + \mathcal{O}(t^2)$$

$$q_{BA}(X_2) = -\frac{htV_3}{2I_{22}}(b - X_2)$$

$$q_{AC}(X_3) = -\frac{htV_3}{2I_{22}}\left(b + \frac{h}{4}\right) + \frac{tV_3}{2I_{22}}X_3^2$$

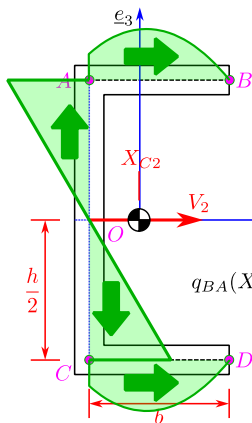
$$q_{CD}(X_2) = -\frac{htV_3}{2I_{22}}(b - X_2)$$

$$M_1 = \oint pqds = -\frac{b^2h^2tV_3}{4I_{22}} := V_3\xi_s$$

$$\xi_s \approx -\frac{3b^2}{h + 6b} + \mathcal{O}(t)$$

## 2. Shear Stress and Flow in Sections

- Shear Center is the point of shear load such that  $\sum m_1 = 0$ . Although in a lot of symmetric sections this is coincident with the centroid, **this is NOT always the case.**
- Consider the “C” section beam:



$$XC_2 = \frac{(bt)\frac{b}{2} \times 2}{2bt + ht} = \frac{b}{2 + h/b}$$

$$q(s) - q(0) = -\frac{tV_2}{I_{33}} \int_0^s X_2 ds$$

### Shear Flow

$$q_{BA}(X_2) = \frac{tV_2}{2I_{33}} \frac{(X_2^2(h+2b)^2 - b^2(h+b)^2)}{(h+2b)^2}$$

$$q_{AC}(X_3) = -\frac{tb^2V_2}{(2b+h)I_{33}} X_3$$

$$q_{CD}(X_2) = -q_{BA}(X_2)$$

## 2. Shear Stress and Flow in Sections

- Choose constant of integration  $q_0$  to exclude torsional effects.

## 2. Shear Stress and Flow in Sections

- If loads at boundary are self equilibrated to begin with, what is the distribution of loads in the member?

### 3. Stringer-Web Idealization

- Jump in shear flow across boom with area  $A_r$ :

$$q_2 - q_1 = -\sigma_{11,1} A_r$$



# References I

- [1] C. T. Sun. **Mechanics of Aircraft Structures**, 2nd edition. Hoboken, N.J: Wiley, June 2006. ISBN: 978-0-471-69966-8 (cit. on p. 2).
- [2] T. H. G. Megson. **Aircraft Structures for Engineering Students**, Elsevier, 2013. ISBN: 978-0-08-096905-3 (cit. on p. 2).